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FORMULA SHEET

for

GATE-AE AERODYNAMICS





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AERODYNAMICS

AERODYNAMICS BASICS

1. Linear Strains of Fluid Element

$$\dot{\epsilon}_{xx} = \frac{\partial u}{\partial x}; \quad \dot{\epsilon}_{yy} = \frac{\partial v}{\partial y}; \quad \dot{\epsilon}_{zz} = \frac{\partial w}{\partial z}$$

where u, v and w are the velocity components in x, y and z directions respectively.

2. Rate of angular deformation:

$$\dot{\gamma}_{xy} = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \quad \dot{\gamma}_{yz} = \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z}$$

$$\dot{\gamma}_{zx} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}$$

3. Vorticity

$$\xi = \nabla \times \vec{V};$$

IF $\nabla \times \vec{V} = 0 \rightarrow$ Irrotational flow

4. Angular Velocity

$$\vec{\omega} = \frac{1}{2} [\nabla \times \vec{V}]$$

5. Enstrophy

$$(E) = |\xi|^2$$

$$(E) = |\nabla \times \vec{V}|^2$$

6. Continuity Equation for

Incompressible Flow:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \Rightarrow \dot{\epsilon}_{xx} + \dot{\epsilon}_{yy} + \dot{\epsilon}_{zz} = 0$$

i.e. No volumetric dilation rate per unit volume

7. Stream Function:

$$\bar{\psi}(x, y) = c$$

For each 'c' value we get a streamline

$$\rho u = \frac{\partial \bar{\psi}}{\partial y}; \quad \rho v = -\frac{\partial \bar{\psi}}{\partial x} \text{ And}$$

$$\rho V_r = \frac{1}{r} \frac{\partial \bar{\psi}}{\partial \theta}; \quad \rho V_\theta = -\frac{\partial \bar{\psi}}{\partial r}$$

For incompressible flow $\psi = \frac{\bar{\psi}}{\rho}$

$$\therefore u = \frac{\partial \psi}{\partial y}; \quad v = -\frac{\partial \psi}{\partial x}$$

Note:

- For compressible flow, $\Delta \bar{\psi} = \bar{\psi}_2 - \bar{\psi}_1$ give the mass flow rate between the two streamlines ψ_1 and ψ_2
- For incompressible flow, $\Delta \psi$ represents volume flow rate per unit depth.
- Stream function is valid for 2D flows only.

8. Velocity Potential Function:

ϕ = Velocity potential function

So, the flow velocity is defined as

$$\vec{V} = \nabla \phi \text{ i.e., } u = \frac{\partial \phi}{\partial x}; \quad v = \frac{\partial \phi}{\partial y}; \quad w = \frac{\partial \phi}{\partial z}$$

Polar:

$$V_r = \frac{\partial \phi}{\partial r}; \quad V_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta}; \quad V_z = \frac{\partial \phi}{\partial z}$$

Spherical:

$$V_r = \frac{\partial \phi}{\partial r}; \quad V_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta}; \quad V_\phi = \frac{1}{r \sin \theta} \left[\frac{\partial \phi}{\partial \phi} \right]$$

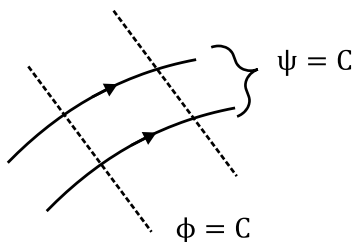
Note:

Velocity potential function is defined for irrotational flow only.

9. Relation b/w stream line and potential lines

$$\left(\frac{dy}{dx}\right)_{\psi=\text{constant}} \times \left(\frac{dy}{dx}\right)_{\phi=\text{constant}} = -1$$

- Streamlines are isolines of stream function.
- Equipotential lines are isolines of potential function.



- Streamlines and Equipotential lines form an orthogonal flow net.

10. Reynolds Transport Theorem

Total derivative of any property

$$\frac{DB_{\text{sys}}}{Dt} = \frac{\partial}{\partial t} \int_{C.V} \rho b dV + \int_{C.S} \rho b (\vec{V} \cdot \vec{n}) dA$$

C.V – Control Volume

C.S. – Control Surface

b- property per unit mass

V-Volume

V- Velocity

A-Control Surface Area

ρ- Density

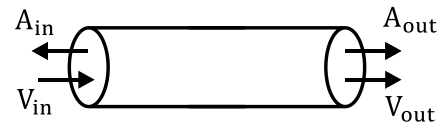
\vec{n} - unit vector normal to control surface

11. Mass continuity Integral form

$$\frac{\partial}{\partial t} \int_{C.V} \rho dV + \int_{C.S} \rho (\vec{V} \cdot \vec{n}) dA = 0$$

12. Momentum Conservation

$$\frac{\partial}{\partial t} \int_{C.V} \vec{V} \rho dV + \int_{C.S} \vec{V} \rho (\vec{V} \cdot \vec{n}) dA = \sum \vec{F} \text{ of C.V}$$



$$\sum \vec{F} = \vec{F}_{\text{pressure}} + \vec{F}_{\text{viscous}} + \vec{F}_{\text{body}} + \vec{F}_{\text{other}}$$

Differential Forms:

13. Continuity Equation:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{V}) = 0$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} (\rho u) + \frac{\partial}{\partial y} (\rho v) + \frac{\partial}{\partial z} (\rho w) = 0$$

Here u, v and w are velocity components.

14. Cylindrical Coordinate Continuity

Equation

$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} (r \rho V_r) + \frac{1}{r} \frac{\partial}{\partial \theta} (\rho V_\theta) + \frac{\partial}{\partial z} (\rho V_z) = 0$$

15. Momentum Conservation

$$\rho \frac{Du}{Dt} = \frac{-\partial p}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + \rho f_x$$

$$\rho \frac{Dv}{Dt} = \frac{-\partial p}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + \rho f_y$$

$$\rho \frac{Dw}{Dt} = \frac{-\partial p}{\partial z} + \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} + \rho f_z$$

Here τ_{ij} - Stress tensor components

p- pressure

f_x, f_y and f_z -body forces

16. Navier-Stokes Equation

Momentum equation for a Newtonian fluid

Stress Terms:

$$\tau_{xx} = \lambda(\nabla \cdot \vec{V}) + 2\mu \frac{\partial u}{\partial x}$$

$$\tau_{yy} = \lambda(\nabla \cdot \vec{v}) + 2\mu \frac{\partial v}{\partial y}$$

$$\tau_{zz} = \lambda(\nabla \cdot \vec{v}) + 2\mu \frac{\partial w}{\partial z}$$

$$\tau_{xy} = \tau_{yx} = \mu \left(\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$

$$\tau_{xz} = \tau_{zx} = \mu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)$$

$$\tau_{yz} = \tau_{zy} = \mu \left(\frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \right)$$

Stokes relation: $\lambda = \frac{-2}{3} \mu$

$\mu \rightarrow$ Molecular viscosity coefficient

$\lambda \rightarrow$ Secondary viscosity coefficient

$\lambda + \frac{2}{3} \mu \rightarrow$ Bulk viscosity

For a Newtonian fluid with constant density and constant viscosity,

$$\rho g_x - \frac{\partial \rho}{\partial x} + \mu \left[\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right] = \rho \frac{Du}{Dt}$$

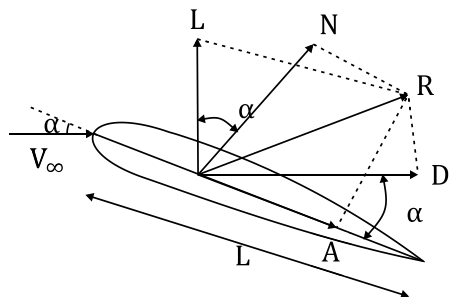
$$\rho g_y - \frac{\partial \rho}{\partial y} + \mu \left[\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right] = \rho \frac{Dv}{Dt}$$

$$\rho g_z - \frac{\partial \rho}{\partial z} + \mu \left[\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right] = \rho \frac{Dw}{Dt}$$

17. Material Derivative

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + (\vec{V} \cdot \nabla)$$

18. Forces on an airfoil:



$$L = N \cos \alpha - A \sin \alpha$$

$$D = N \sin \alpha + A \cos \alpha$$

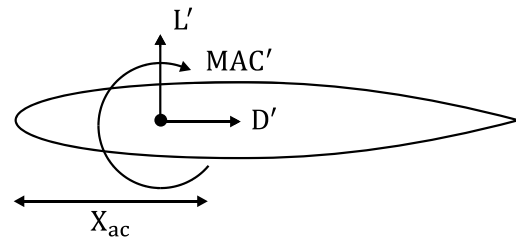
$N \rightarrow$ Normal force; $A \rightarrow$ Axial Force

$L \rightarrow$ Lift force; $D \rightarrow$ Drag Force

$\alpha \rightarrow$ Angle of attack

$V_\infty \rightarrow$ free stream velocity

19. Aerodynamic Center:



$$\bar{x}_{ac} = \frac{-m_0}{a_0} + 0.25$$

where, $a_0 = \frac{dC_L}{d\alpha}$ and

$$m_0 = \frac{dC_{m,c/4}}{d\alpha}$$

20. Aerodynamic Coefficients:

Lift coefficient,

$$C_L = \frac{L}{q_\infty S} ; C_N = \frac{N}{q_\infty S}$$

Drag coefficient,

$$C_D = \frac{D}{q_\infty S} ; C_A = \frac{A}{q_\infty S}$$

Moment Coefficient,

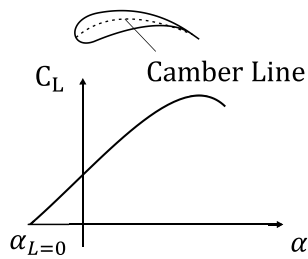
$$C_M = \frac{M}{q_\infty S \bar{c}}$$

Note: $C_L = f_n(\alpha, Re, Mach \text{ no.})$

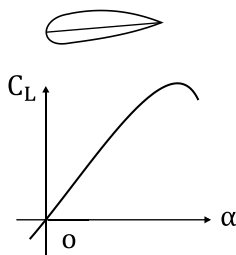
21. Coefficient of lift v/s angle of attack

C_L v/s α :

Cambered airfoil



Symmetric airfoil:



22. NACA Nomenclature:

Four digits: NACA 2412

First Digit: Maximum camber in hundredths of chord, Ex: Maximum camber 0.02c

Second Digit: Location of maximum camber along the chord from the leading edge in tenths of chord. Ex: 0.4c here

Last 2 digits: Maximum thickness in hundredths of chord. Ex: 0.12c

Five Digit: NACA 23012

First Digit: multiplied by $3/2$ gives the design lift coefficient in tenths.

Ex: $\frac{2 \times 3/2}{10} = 0.3$

Second and Third Digit: Divided by 2 gives the location of maximum camber along the chord from the leading edge in hundredths of chord.

Ex: $\frac{30/2}{100} c = 0.15c$

Last 2 Digits: Maximum thickness in hundredths of chord

Ex: $\frac{12}{100} c = 0.12c$

Six-Digit Series: NACA 65-218

First Digit: Identifies series [6]

Second Digit: Location of minimum pressure [0.5c] in tenths of chord from leading edge.

Third Digit: Design lift coefficient in tenths; $C_L = 0.2$

Last 2 Digits: Maximum thickness in hundredths of chord
0.18c or 18% thickness

23. Pressure Coefficient:

$$C_p = \frac{p - p_\infty}{q_\infty}$$

q_∞ – Dynamic pressure

[Applicable from incompressible to hypersonic flow]

$$C_p = 1 - \left(\frac{V}{V_\infty}\right)^2$$

[For incompressible flow only]

Note: C_p is non-dimensional quantity and if flow conditions are similar (for kinematic similarity), we will get the same C_p distribution.

POTENTIAL FLUID FLOW

1. Assumptions:

a. Flow is inviscid and irrotational, thereby no vorticity.

$$\vec{\xi} = \nabla \times \vec{V} = 0$$

Note:

- Potential function was chosen as to satisfy irrotationality condition.
- Similarly stream function was chosen to satisfy continuity equation.

2. Potential function and Laplace

Equation

$$\nabla \cdot \vec{V} = 0$$

By mathematics, \vec{V} should be gradient of scalar function, then,

$$\nabla \cdot (\nabla \phi) = 0 \text{ or } \nabla^2 \phi = 0$$

ϕ satisfies Laplace equation, then only flow is possible as it satisfies continuity equation.

Similarly, $\nabla^2 \psi = 0$ for stream function

Note:

- Any irrotational, incompressible flow has a velocity potential and stream function (for 2D) that both satisfy Laplace equation.
- Conversely, any solution of Laplace equation represents the velocity potential or stream function (2D) for an irrotational, incompressible flow.

3. Features of Laplace Equation:

$$\phi = \phi_1 + \phi_2 + \phi_3 + \dots + \phi_n$$

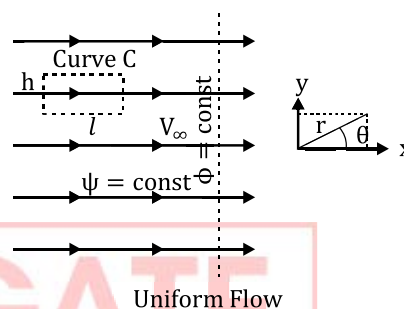
- Combining several potential flows gives us a new potential flow.
- Each of the potential flow is called an elementary flow.

- We can add their derivatives too, i.e., velocities (Observe elementary flows combined).

4. Circulation:

$$\Gamma = - \oint V ds = - \iint_A (\nabla \times V) dA$$

5. Uniform Flow:



In x direction, V_∞ – velocity

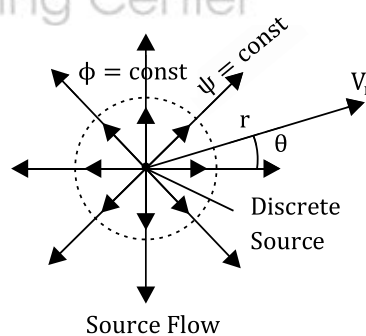
Potential function,

$$\phi = V_\infty x = V_\infty r \cos \theta$$

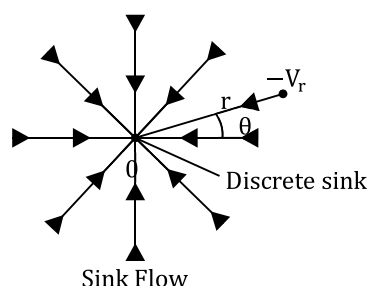
$$\text{Stream function, } \psi = V_\infty y = V_\infty r \sin \theta$$

circulation taken in the loop C, $\Gamma = 0$

6. Source (\wedge) and sink ($-\wedge$) Flow:



Source Flow



Sink Flow

Radial velocity, $V_r = \frac{\Lambda}{2\pi r}$,

Tangential $V_\theta = 0$

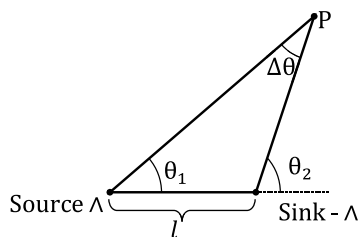
Source strength $\Lambda = 2\pi r V_r \text{ m}^2/\text{s}$

Potential and stream function

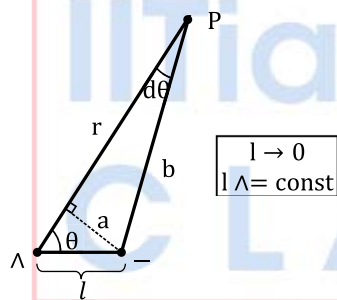
$$\phi = \frac{\Lambda}{2\pi} \ln r ; \quad \psi = \frac{\Lambda}{2\pi} \theta ;$$

Circulation, $\Gamma = 0$

7. Doublet Flow:



(a) Source - sink pair



(b) Limiting case for a doublet

Figure: How a source-sink pair

approaches a doublet in the limiting case.

Doublet Strength $k = \Lambda l$ and $\Gamma = 0$

Potential and stream function

$$\psi = \frac{-k \sin \theta}{2\pi r} ; \quad \phi = \frac{k \cos \theta}{2\pi r}$$

Stream line

$\psi = \text{constant}$

$$= -\frac{k \sin \theta}{2\pi r} = C$$

$$r = -\frac{k}{2\pi C} \sin \theta$$

$$r = d \sin \theta$$

d-dia of Lobe

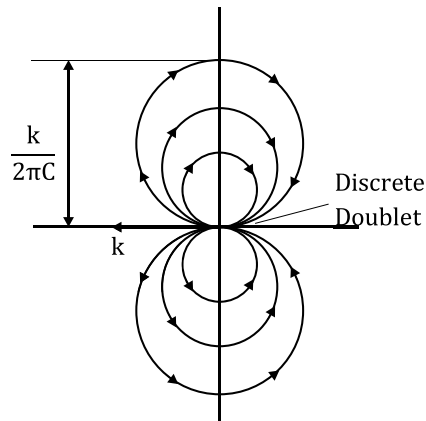


Figure: Doublet Flow with Strength k.

8. Vortex Flow:

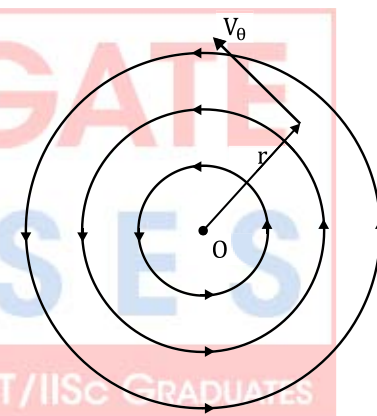


Figure: Vortex Flow.

Vortex Strength is Γ (Circulation)

$$V_\theta = -\frac{\Gamma}{2\pi r} ; \quad V_r = 0$$

$$\Gamma = -2\pi r = \text{Constant}$$

Potential and stream function

$$\phi = \frac{-\Gamma}{2\pi} \theta ; \quad \psi = \frac{\Gamma}{2\pi} \ln r$$

A vortex of positive strength (Γ) rotates in clockwise direction.

Types of flow	Velocity	ϕ	ψ
Uniform flow in x-direction	$u = V_\infty; v = 0$	$V_\infty x$	$V_\infty y$
Source flow	$V_r = \frac{\Lambda}{2\pi r}; V_\theta = 0$	$\frac{\Lambda}{2\pi} \ln r$	$\frac{\Lambda}{2\pi} \theta$
Doublet Flow	$V_r = \frac{-k \cos \theta}{2\pi r^2}; V_\theta = \frac{-k \sin \theta}{2\pi r^2}$	$\frac{k \cos \theta}{2\pi r}$	$\frac{-k \sin \theta}{2\pi r}$
Vortex Flow	$V_r = 0; V_\theta = -\frac{\Gamma}{2\pi r}$	$\frac{-\Gamma}{2\pi} \theta$	$\frac{\Gamma}{2\pi} \ln r$

Source + Uniform Flow:

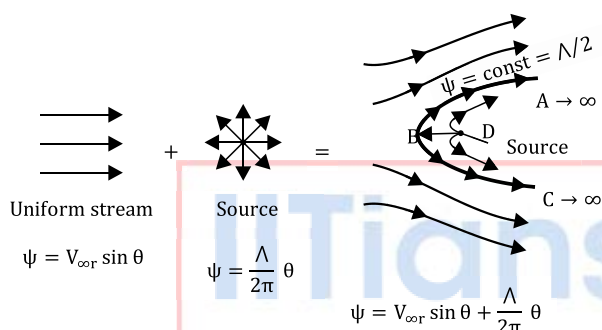


Figure: Superposition of a uniform flow and a source; flow over a semi-infinite body.

$$V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} \text{ and } V_\theta = -\frac{\partial \psi}{\partial r}$$

→ Applicable for all polar coordinates

Stagnation points,

$$(r, \theta) = \left(\frac{\Lambda}{2\pi V_\infty}, \pi \right)$$

Stagnation streamline,

$$\psi = \frac{\Lambda}{2} = \text{constant}$$

Width of Rankine half body

$$b = \frac{\Lambda}{V_\infty}$$

9. Source + Sink + Uniform Flow

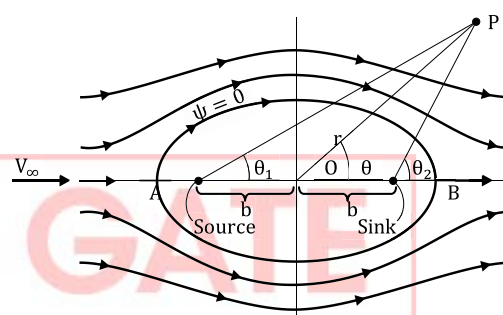


Figure: Superposition of a uniform flow and a source-sink pair; flow over a Rankine oval

$$\psi = V_\infty r \sin \theta + \frac{\Lambda}{2\pi} \theta_1 - \frac{\Lambda}{2\pi} \theta_2$$

Stagnation Points

$$OA = OB = \sqrt{b^2 + \frac{\Lambda b}{\pi V_\infty}}$$

and are at $\theta = 0, \pi$

Stagnation Streamline, $\psi = 0$

10. Non-Lifting Flow Over Cylinder

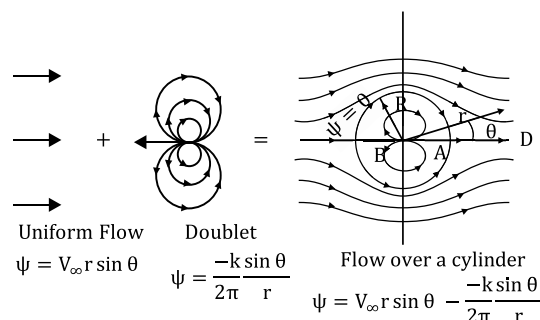


Figure: Superposition of a uniform flow and a doublet; non-lifting flow over a circular cylinder.

$$\psi = V_{\infty} r \sin \theta \left(1 - \frac{R^2}{r^2} \right)$$

Where, $R^2 = \frac{k}{2\pi V_{\infty}}$

$$\begin{cases} V_r = V_{\infty} \cos \theta \left(1 - \frac{R^2}{r^2} \right) \\ V_{\theta} = -V_{\infty} \sin \theta \left(1 + \frac{R^2}{r^2} \right) \end{cases}$$

There are two stagnation points, $(R, 0)$ and (R, π)

At stagnation streamline, $\psi = 0$

On cylinder surface, at all points $r = R$

$$V_r = 0; V_{\theta} = -2V_{\infty} \sin \theta$$

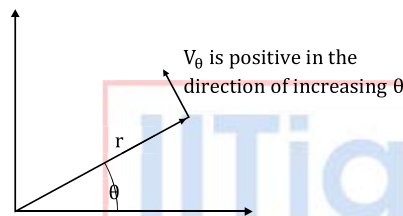
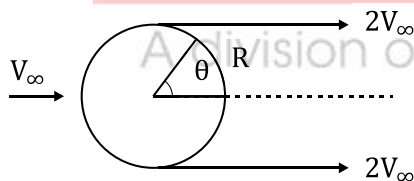


Figure: Sign Convention for V_{θ} in polar coordinates.

Note:

1. The maximum velocity is $2V_{\infty}$ at top and bottom of the cylinder.



Maximum acceleration [Flow $r \rightarrow R$] occurs

at $\theta = 135^\circ, 225^\circ$ and $a_{\max} = \frac{2V_{\infty}^2}{R}$

Maximum deceleration occurs at $\theta = 45^\circ, 315^\circ$

2. Pressure coefficient

$$C_p = 1 - \left(\frac{V}{V_{\infty}} \right)^2$$

$$C_p = 1 - 4 \sin^2 \theta$$

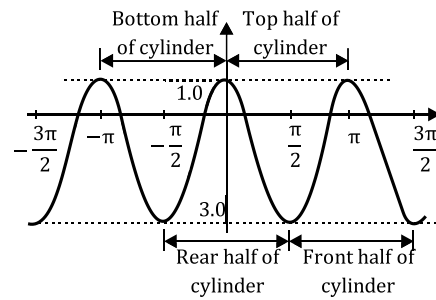


Figure: Pressure coefficient distribution over the surface of a circular cylinder, theoretical results for inviscid, incompressible flow.

11. Lifting Flow Over Cylinder:

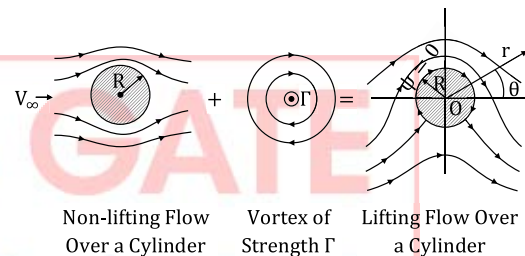


Figure: The synthesis of lifting flow over a circular cylinder.

Non-lifting flow over cylinder + Vortex flow
Stream function

$$\psi = (V_{\infty} r \sin \theta) \left(1 - \frac{R^2}{r^2} \right) + \frac{\Gamma}{2\pi} \ln \frac{r}{R}$$

V_r and V_{θ} can be achieved simply by adding velocities of non-lifting flow and vortex flow.

$$V_r = \left(1 - \frac{R^2}{r^2} \right) V_{\infty} \cos \theta$$

$$V_r = - \left(1 + \frac{R^2}{r^2} \right) V_{\infty} \sin \theta - \frac{\Gamma}{2\pi r}$$

stagnation Point $r = R$

$$\theta = \sin^{-1} \left[\frac{-\Gamma}{4\pi V_{\infty} R} \right]$$

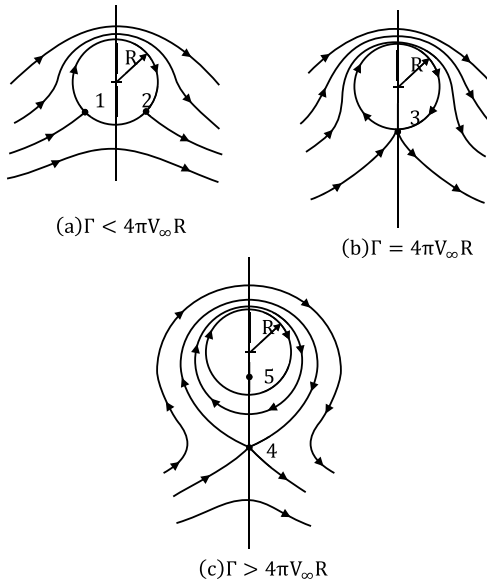


Figure: Stagnation points for the lifting flow over a circular cylinder.

- $\Gamma < 4\pi V_\infty R \rightarrow 2$ stagnation points
- $\Gamma = 4\pi V_\infty R \rightarrow 1$ stagnation point $\rightarrow @ [R, -\frac{\pi}{2}]$
- $\Gamma > 4\pi V_\infty R \rightarrow 1$ stagnation point inside cylinder and 1 stagnation point lifts off the surface of cylinder.

And this condition is also satisfied by $r =$

$$R \text{ and } \theta = \pi/2 \text{ or } -\pi/2$$

So,

$$r = \frac{\Gamma}{4\pi V_\infty} \pm \sqrt{\left(\frac{\Gamma}{4\pi V_\infty}\right)^2 - R^2}$$

Velocity on the surface of cylinder:

$$V = V_\theta = -2V_\infty \sin \theta - \frac{\Gamma}{2\pi R}$$

12. Kutta-Joukowski Theorem:

Lift per unit span,

$$L' = \rho_\infty V_\infty \Gamma$$

$$C_L = \frac{\Gamma}{RV_\infty}$$

■ INCOMPRESSIBLE FLOW OVER AIRFOILS

- Air foil leading edge is usually circular with radius of approx. $0.02c$ and c is chord length.
- Higher the C_{Lmax} , lower is the stalling speed as $(L = \frac{1}{2} \rho_\infty V_\infty^2 SC_L)$
- $\alpha_{L=0}$ is negative for positively cambered airfoils and positive for negative cambered airfoils.
- Lift Slope a_0 is not influenced by Re , but C_{Lmax} increases with increase in Re
- The moment coefficient is also insensitive to Re except at large α
- C_d is sensitive to Re and decreases at higher Re for the order of $Re \sim 10^6$.
- Aerodynamic center, $C_{m,ac} = \text{constant}$

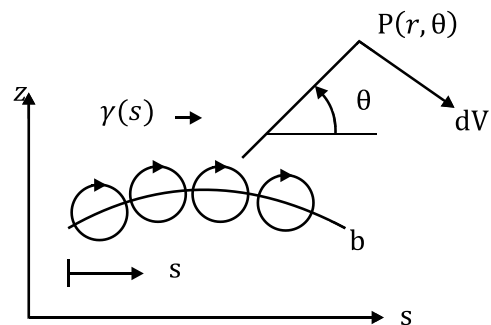
$$\frac{\partial C_m}{\partial \alpha} \bigg|_{ac} = 0$$

1. Vortex Sheet:

$$\phi(s, z) = \frac{-1}{2\pi} \int_a^b \theta \gamma ds$$

$$\gamma = \int_a^b \gamma(s) \cdot ds \quad ; \quad dV = \frac{-\gamma ds}{2\pi r}$$

$$\gamma(s) = \frac{d\Gamma}{ds}$$



2. **The local jump in tangential velocity across the vortex sheet is equal to local sheet strength.**

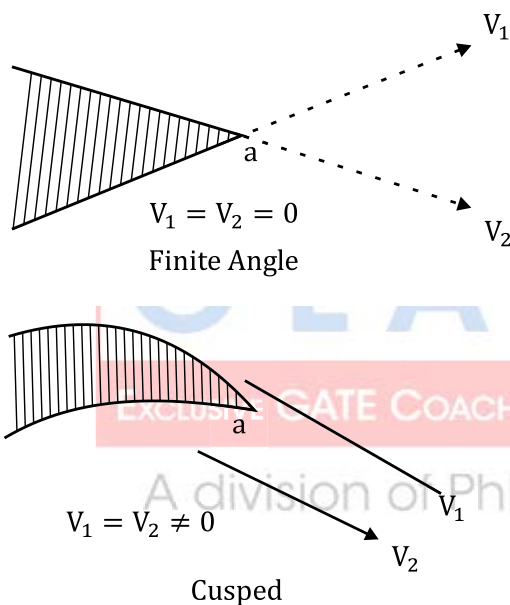
$$\gamma = u_1 - u_2$$

Note: The strength of vortex sheet $\gamma(s)$ is calculated such that, the camber line become a streamline of the flow.

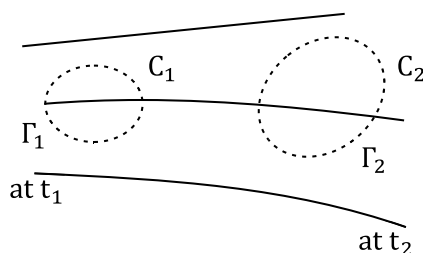
3. **Kutta Condition:**

$$\gamma(T.E) = 0$$

Trailing edge is a stagnation point for finite angled TE airfoils.



4. **Kelvin's Circulation Theorem:**



$$\Gamma_1 = \Gamma_2$$

$$\text{Total derivative } \frac{D\Gamma}{Dt} = 0$$

The time rate of change of circulation around a closed curve consisting of same fluid elements is zero.

THIN AIRFOIL THEORY

For the camber line to be a streamline of flow. Net normal velocity must be zero.

$$V_{\infty, n} + w(s) = 0$$

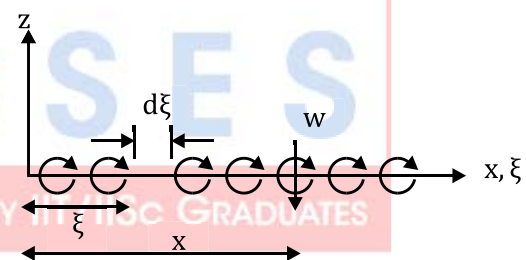
$$V_{\infty, n} = V_{\infty} \sin \left[\alpha + \tan^{-1} \left(\frac{-dz}{dx} \right) \right]$$

$$\text{Approximating, i.e., } V_{\infty, n} = V_{\infty} \left(\alpha - \frac{dz}{dx} \right)$$

$$\text{and } w(s) \approx w(x); \quad w(x) = - \int_0^c \frac{\gamma(\xi) d\xi}{2\pi(x - \xi)}$$

Fundamental equation of thin airfoil theory

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{(x - \xi)} = V_{\infty} \left(\alpha - \frac{dz}{dx} \right)$$



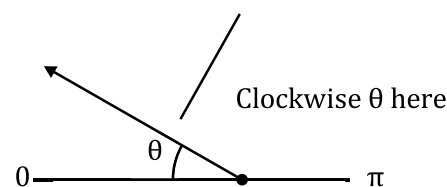
Note:

For symmetric airfoil camber line coincides with chord line,

$$\text{i.e., } \frac{dz}{dx} = 0$$

Symmetric Airfoil

$$\gamma(\theta) = 2\alpha V_{\infty} \frac{(1 + \cos \theta)}{\sin \theta}$$



$$\bullet \quad C_L = 2\pi\alpha$$

- Lift slope = $\frac{dC_L}{d\alpha} = 2\pi$

- Moment coefficient.

$$C_{m,le} = -\frac{C_L}{4}$$

$$C_{m,c/4} = C_{m,ac} = 0$$

- The aerodynamic center and center of pressure are both at the quarter chord point.

Cambered Airfoil: Circulation

$$\Gamma = cV_\infty \left[\pi A_0 + \frac{\pi}{2} A_1 \right]$$

$$\gamma(\theta) = 2V_\infty \left[A_0 \left(\frac{1 + \cos \theta}{\sin \theta} \right) + \sum_{n=1}^{\infty} A_n \sin \theta \right]$$

$$C_L = 2\pi \left[\alpha + \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos \theta_0 - 1) d\theta_0 \right]$$

$$C_L = \pi(2A_0 + A_1)$$

$$\alpha_{L=0} = \frac{-1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos \theta_0 - 1) d\theta_0$$

$$[\because C_L = 2\pi(\alpha - \alpha_{L=0})]$$

$$C_{m,le} = -\frac{C_L}{4} + \frac{\pi}{4} (A_1 - A_2)$$

$$C_{m,c/4} = \frac{\pi}{4} (A_2 - A_1)$$

Usually, $A_1 > A_2$

$$X_{cp} = \frac{c}{4} \left[1 + \frac{\pi}{C_L} (A_1 - A_2) \right]$$

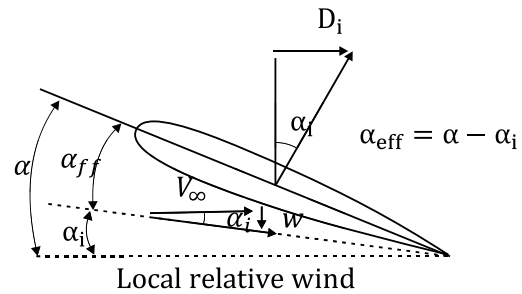
Fourier Coefficients:

$$A_0 = \alpha - \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} d\theta_0$$

$$A_n = \frac{2}{\pi} \int_0^\pi \frac{dz}{dx} \cos n\theta_0 d\theta_0$$

- Lift slope = $\frac{dC_L}{d\alpha} = 2\pi$
- The aerodynamics center is at the quarter-chord point.
- The center of pressure varies with the lift coefficient.

INCOMPRESSIBLE FLOW OVER WING



$\omega \rightarrow$ downwash.

D_i induced drag $[L \sin \alpha_1]$

For airfoil, $C_l = 2\pi\alpha$;

For wing $C_L = a(\alpha - \alpha_i)$

1. Definitions:

Geometric twist: α is different at different spanwise locations.

Washout: If the tip is at a higher α than the root

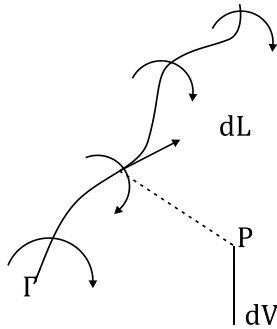
Aerodynamic Twist: Different airfoil sections along the span with different values of $\alpha_{L=0}$

2. Biot Savart Law:

The strength of the vortex

filament is defined as Γ . Consider a directed segment of the filament dl , as shown in Figure. The radius vector from dl to an arbitrary point P in space is r . The segment dl induces a velocity at P

$$dV = \frac{\Gamma}{4\pi} \frac{dL \times r}{|r|^3}$$



3. Velocity induced at a point 'P' by an infinite straight vortex filament.

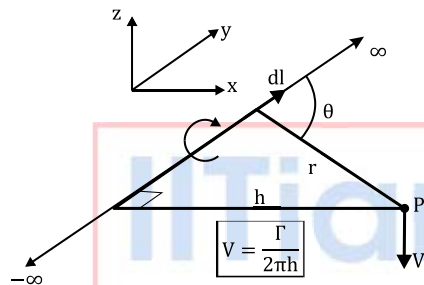


Figure: Velocity induced at point P by an infinite, straight vortex filament.

$$V = \frac{\Gamma}{2\pi h}$$

4. Semi-infinite Vortex:

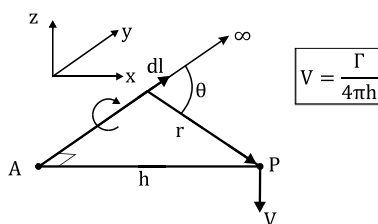


Figure: Velocity induced at point P by a semi-infinite straight vortex filament.

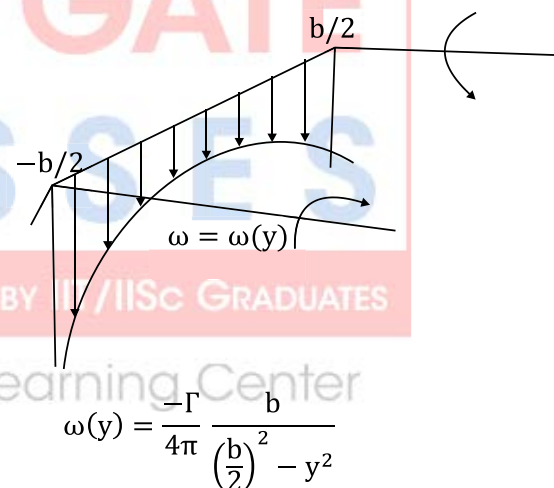
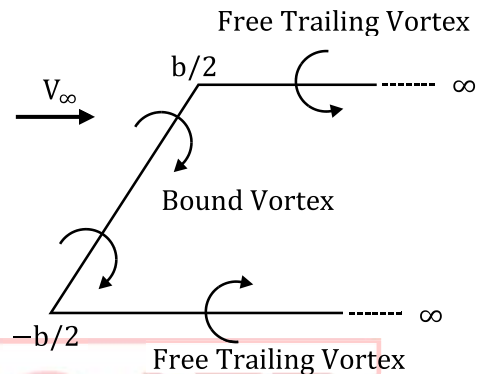
$$V = \frac{\Gamma}{4\pi h}$$

5. Helmholtz's Vortex Theorems:

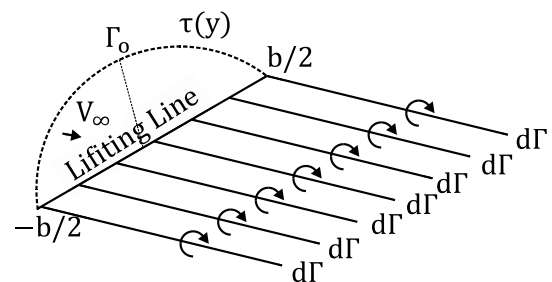
- i. The strength of a vortex filament is constant along its length.

- ii. A vortex filament cannot end in a fluid it must extend to the boundary of the fluid (which can be $\pm\infty$) or form a closed path.

6. Prandtl's lifting line theory



Note: At trailing edge $y = \pm \frac{b}{2}$, downwash is infinity. So, this model was not accurate.



$$\omega(y_0) = \frac{-1}{4\pi} \int_{-\frac{b}{2}}^{\frac{b}{2}} \left(\frac{d\Gamma}{dy} \right) \frac{dy}{y_0 - y}$$

$$\alpha_i(y_0) = \tan^{-1} \left(\frac{-\omega(y_0)}{V_\infty} \right) = \alpha - \frac{\omega(y_0)}{V_\infty}$$

$$\therefore \alpha_i(y_0) = \frac{1}{4\pi V_\infty} \int_{-\frac{b}{2}}^{\frac{b}{2}} \left(\frac{d\Gamma}{dy} \right) \frac{dy}{y_0 - y}$$

Elliptical Lift Distribution:

$$\Gamma(y) = \Gamma_0 \sqrt{1 - \left(\frac{2y}{b} \right)^2}$$

\therefore Circulation and Lift distribution

$$L'(y) = \rho_\infty V_\infty \Gamma_0 \sqrt{1 - \left(\frac{2y}{b} \right)^2}$$

$$\text{Downwash, } \omega(\theta_0) = -\frac{\Gamma_0}{2b}$$

\rightarrow Constant downwash for elliptic distribution

Induced angle of attack,

$$\alpha_i = \frac{\Gamma_0}{2bV_\infty}; \left[\alpha_i = \frac{\omega}{V_\infty} \right]$$

$$\text{Lift, } L = \rho_\infty V_\infty \Gamma_0 \frac{b}{4} \pi$$

Alternative:

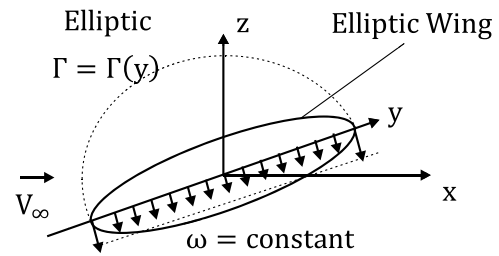
Induced angle of Attack:

$$\alpha_i = \frac{C_L}{\pi AR}$$

$$\text{Induced drag, } C_{Di} = \frac{C_L^2}{\pi AR}$$

AR- Aspect Ratio

Note: For elliptic lift distribution the wing platform is elliptical



General Lift Distribution:

$$C_{Di} = \frac{C_L^2}{\pi AR} (1 + \delta) = \frac{C_L^2}{\pi e AR}$$

$$\frac{1}{1 + \delta} = e$$

\rightarrow Ostwalds Span efficiency factor

Taper ratio = $\frac{c_t}{c_r}$; $c_t \rightarrow$ tip chord;
 $c_r \rightarrow$ root chord

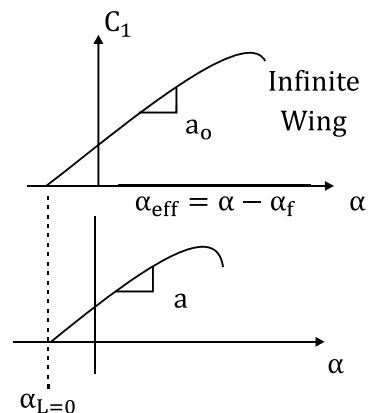
Lift Curve Slope:

$$\text{For airfoil, } a_0 = \frac{dC_L}{d\alpha}$$

$$\text{Finite wing, } a = \frac{dC_L}{d\alpha}$$

$$a_0 > a$$

$$a = \frac{a_0}{1 + \frac{a_0}{\pi e AR}} \text{ (in radian)}$$



Note:

- a. If we use α_{eff} ($\alpha - \alpha_i$) and calculate C_L using slope, a_0 then,
$$C_L = a_0(\alpha_{\text{eff}} - \alpha_{L=0})$$
- b. If we use geometric AOA α , then slope used is a
$$C_L = a(\alpha - \alpha_{L=0})$$



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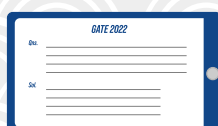
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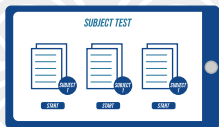
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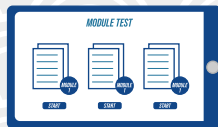
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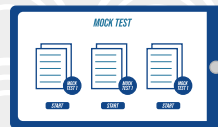
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